

2nd semester CC-03 Biot-Savart's law

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The Biot Savart Law is used to determine the magnetic field intensity H near a current-carrying conductor or we can say, it gives the relation between magnetic field intensity generated by its source current element. The law was stated in the year 1820 by Jean Baptiste Biot and Felix Savart. The direction of the magnetic field follows the right hand rule for the straight wire. Biot Savart law is also known as Laplace's law or Ampere's law.

Biot Savart Law states that the magnetic field at any point due to a steady current in an infinitely long straight wire is directly proportional to the current and inversely proportional to the distance from point to wire.

Any current carrying element is the source of magnetic fields. When charges move in a conducting wire and produce a current I , the magnetic field $d\vec{B}$ due to an infinitesimal current element $I d\vec{l}$ at a point P is given by

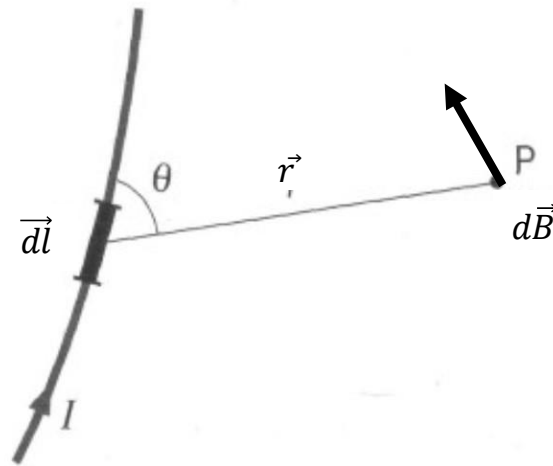
$$d\vec{B} = \frac{\mu_0 I d\vec{l} \times \vec{r}}{4\pi r^3} \quad (1)$$

where \vec{r} is the line vector joining the current element and the point P, and μ_0 is a constant called the permeability of free space

$$\mu_0 = 4\pi \times 10^{-7} \text{ T.m/A}$$

Direction of $d\vec{B}$ is perpendicular to the plane perpendicular to the plane containing both the $d\vec{l}$ and \vec{r} following right hand rule.

To apply the equation, the point in space where the magnetic field is to be calculated is arbitrarily chosen r . Holding that point fixed, the line integral over the path of the electric current is calculated to find the total magnetic field at that point. The application of this law implicitly relies on the superposition principle for magnetic fields, i.e. the fact that the magnetic field is a vector sum of the field created by each infinitesimal section of the wire individually.



The total magnetic field \vec{B} at point P due to the current carrying conductor can be calculated by integrating adding up the magnetic field contributions from segments

$$\vec{B} = \int_C \frac{d\vec{B}}{c} = \frac{\mu_0}{4\pi c} \int_C \frac{I d\vec{l} \times \vec{r}}{r^3} \quad (2)$$

The integral is a vector integral, which means that the expression for \vec{B} is really three integrals, one for each component of \vec{B} .

The integral is usually around a closed curve, since stationary electric currents can only flow around closed paths when they are bounded. However, the law also applies to infinitely long wires.

The Lorentz force law

In SI units, the force \vec{F} acting on a particle of electric charge q with instantaneous velocity \vec{v} , due to an external electric field \vec{E} and magnetic field \vec{B} , is given by

$$\vec{F} = q(\vec{E} + \vec{v} \times \vec{B}) \quad (3)$$

The term $q\vec{E}$ is called the electric force, while the term $q\vec{v} \times \vec{B}$ is called the magnetic force. According to some definitions, the term "Lorentz force" refers specifically to the formula for the magnetic force, with the total electromagnetic force (including the electric force) given some other (nonstandard) name. In general, the term "Lorentz force" refers to the expression for the total force.

For a continuous charged distribution in motion, the Lorentz force equation becomes

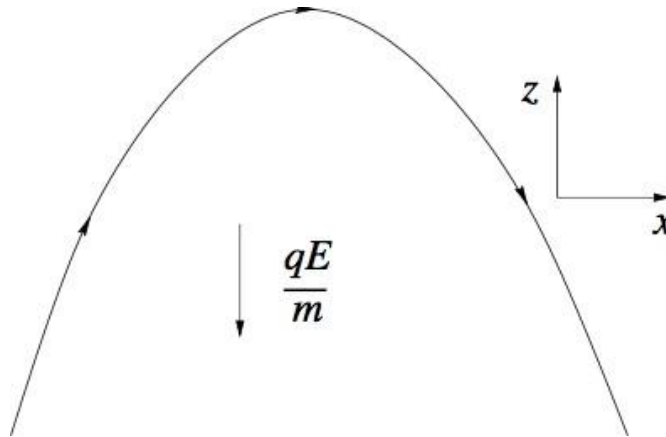
$$d\vec{F} = d(E + \vec{v} \times \vec{B}) \quad (4)$$

where $d\vec{F}$ is the force on a small piece of the charged distribution with charge dq .

Motion of the particle under the Lorentz force Electric field (electric) alone

$$m \frac{d\vec{v}}{dt} = q\vec{E} \quad (5)$$

Orbit depends only on ratio q/m . For uniform field this equation is analogous to particle moving under influence of gravity. Acceleration $g \leftrightarrow \frac{qE}{m}$. Orbits are parabolas.



Energy is conserved taking into account potential energy $q\phi$. [**Homework: prove it!**] **Magnetic field (uniform) alone**

$$\frac{d\vec{v}}{dt} = q\vec{v} \times \vec{B}$$

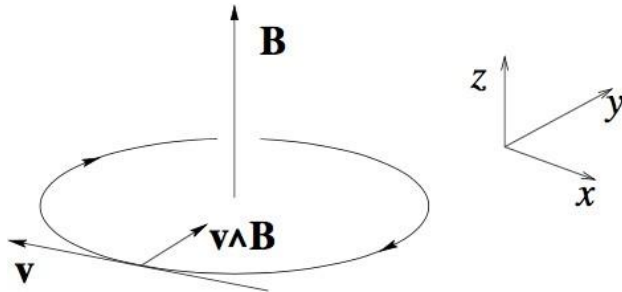
Taking \vec{B} in \hat{z} direction, z -component of the velocity v_z is constant.

$$\begin{aligned} \dot{v}_x &= \frac{q}{m} v_y B & \dot{v}_y &= \frac{-q}{m} v_x B \\ \Rightarrow \ddot{v}_x &= -\left(\frac{qB}{m}\right)^2 v_x & \ddot{v}_y &= -\left(\frac{qB}{m}\right)^2 v_y \end{aligned} \quad (6)$$

$$\begin{aligned} v_x &= v \sin \frac{qB}{m} t & v_y &= v \cos \frac{qB}{m} t \\ x &= -v \frac{m}{qB} \cos \frac{qB}{m} t + x_0 & y &= v \frac{m}{qB} \sin \frac{qB}{m} t + y_0 \end{aligned}$$

The equation of a circle with center at (x, y) and radius $\left(\frac{vm}{qB}\right)$.

qB



Helical motion results when the velocity vector is not perpendicular to the magnetic field vector i.e. having a constant non-zero z component.

The motion that is produced when one component of the velocity is constant in magnitude and direction (i.e., straight-line motion) while the other component is constant in speed but uniformly varies in direction (i.e., circular motion). It is the superposition of straight-line and circular motion.

Cycloid motion

The force on the charged particle in the presence of both electric and magnetic fields is given

$$m \frac{d\vec{v}}{dt} = q\vec{v} \times \vec{B} + q\vec{E}$$

Let the \vec{E} and \vec{B} be at right angle to each other, so that $\vec{E} \cdot \vec{B} = 0$

If the particle is initially at rest no magnetic force acts on the particle. As the electric field exerts a force on the particle, it acquires a velocity in the direction of \vec{E} . The magnetic force now acts sideways on the particle.

For a quantitative analysis of the motion, let \vec{E} be taken along the x -direction and \vec{B} along z -direction. As there is no component of the force along the z -direction, the velocity of the particle remains zero in this direction. The motion, therefore, takes place in $x-y$ plane. The equations of motion are

$$m \frac{dv_x}{dt} = (E + v_y B) = qE + qBv_y \quad (7)$$

$$m \frac{dv_y}{dt} = q(E + v_x B) \rightarrow -qBv_x \quad (8)$$

Differentiating equation 7 and putting the value from equation (8)

$$m \frac{d^2 v_x}{dt^2} = qB \frac{dv_y}{dt} = -q \frac{dB}{m} v_x \quad (9)$$

which has the solution

$$v_x = A \sin \omega_c t \quad (10)$$

where $\omega_c = \frac{qB}{m}$

Substituting this solution into the equation for v_y , we get

$$v_y = a \cos \omega_c t + C \quad (11)$$

The constant A may be determined by substituting the solutions in eqn. (7) which gives

$$m \omega_c A \cos \omega_c t = qE + qB(\cos \omega_c t - 1) \quad (12)$$

Since the equation above is valid for all times, the constant terms on the right must cancel, which gives $A = E/B$. Thus we have

$$v_x = \frac{E}{B} \sin \omega_c t \quad (13)$$

$$v_y = \left(\frac{E}{B} \cos \omega_c t - 1 \right) \quad (14)$$

The equation to the trajectory is obtained by integrating the above equation and determining the constant of integration from the initial position (taken to be at the origin),

$$x = \frac{E}{B\omega_c} (1 - \cos \omega_c t) \quad (15)$$

$$y = \frac{E}{B\omega_c} (\sin \omega_c t - \omega_c t) \quad (16)$$

The equation to the trajectory is

$$\left(x - \frac{E}{B\omega_c} \right)^2 + \left(y - \frac{Et}{B\omega_c} \right)^2 = \frac{E^2}{B^2\omega_c^2}$$

which represents a circle of radius $R = \frac{E}{B\omega_c}$

whose centre travels along the negative y direction with a constant speed $v = \frac{E}{B}$

The trajectory resembles that of a point on the circumference of a wheel of radius R , rolling down the y-axis without slipping with a speed v . The trajectory is known as a cycloid.