



Angular Momentum

We know from our previous discussion that the commutator \hat{A} and \hat{B} , is defined as

$$[\hat{A}, \hat{B}] \equiv \hat{A}\hat{B} - \hat{B}\hat{A}$$

The following commutator identities are helpful in evaluating commutators:

$$[\hat{A}, \hat{B}] = -[\hat{B}, \hat{A}] \dots\dots\dots(1)$$

$$[\hat{A}, \hat{A}^n] = 0, n=1,2,3\dots \dots\dots(2)$$

$$[k\hat{A}, \hat{B}] = [\hat{A}, k\hat{B}] = k[\hat{A}, \hat{B}] \dots\dots\dots(3)$$

$$[\hat{A}, \hat{B} + \hat{C}] = [\hat{A}, \hat{B}] + [\hat{A}, \hat{C}], \quad [\hat{A} + \hat{B}, \hat{C}] = [\hat{A}, \hat{C}] + [\hat{B}, \hat{C}] \dots\dots\dots(4)$$

$$[\hat{A}, \hat{B}\hat{C}] = [\hat{A}\hat{B}]\hat{C} + \hat{B}[\hat{A}, \hat{C}], \quad [\hat{A}\hat{B}, \hat{C}] = [\hat{A}, \hat{C}]\hat{B} + \hat{A}[\hat{B}, \hat{C}] \dots\dots\dots(5)$$

Where k is constant and the operator are assumed to be linear.

Example:

1. $[\partial/\partial x, x] = 1$

Proof: $[\partial/\partial x, x]f(x) = \partial/\partial x \{x.f(x)\} - x\partial/\partial x \{f(x)\}$

$$= xf'(x) + f(x) - xf'(x) = 1.f(x)$$

So, $[\partial/\partial x, x]f(x) = 1.f(x)$

Or, $[\partial/\partial x, x] = 1$ (proved) $\dots\dots\dots(6)$

2.

$$[\hat{x}, \hat{p}_x] = i\hbar$$

Proof: $[\hat{x}, \hat{p}_x] = [x, \frac{\hbar}{i} \partial/\partial x] = \frac{\hbar}{i} [x, \partial/\partial x]$

$$= -\frac{\hbar}{i} [\partial/\partial x, x] = -\frac{\hbar}{i}$$

$$[\hat{x}, \hat{p}_x] = i\hbar$$

$\dots\dots\dots(7)$

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3.

$$\begin{aligned}
 [\hat{x}, \hat{p}_x^2] &= 2\hbar^2 \partial/\partial x \\
 \text{Proof: } [\hat{x}, \hat{p}_x^2] &= [\hat{x}, \hat{p}_x] \hat{p}_x + \hat{p}_x [\hat{x}, \hat{p}_x] \\
 &= i\hbar \frac{\hbar}{i} \partial/\partial x + \frac{\hbar}{i} \partial/\partial x i\hbar \\
 &= 2\hbar^2 \partial/\partial x
 \end{aligned}
 \dots\dots\dots(8)$$

4.

$$\begin{aligned}
 [\hat{x}, \hat{H}] &= [\hat{x}, \hat{T} + \hat{V}] = [\hat{x}, \hat{T}] + [\hat{x}, \hat{V}(x,y,z)] = [\hat{x}, \hat{T}] \\
 &= [x, (1/2m)(\hat{p}_x^2 + \hat{p}_y^2 + \hat{p}_z^2)] \\
 &= (1/2m)[\hat{x}, \hat{p}_x^2] + (1/2m)[\hat{x}, \hat{p}_y^2] + (1/2m)[\hat{x}, \hat{p}_z^2] \\
 &= (1/2m) \cdot 2\hbar^2 \partial/\partial x + 0 + 0 \\
 [\hat{x}, \hat{H}] &= \frac{\hbar^2}{2m} \partial/\partial x = \frac{i\hbar}{m} \hat{p}_x \quad (\text{Proved})
 \end{aligned}
 \dots\dots\dots(9)$$

The above commutators have important physical consequences . Since $[\hat{x}, \hat{p}_x] \neq 0$, we cannot expect the state function to be simultaneously an eigenfunction of \hat{x} and of \hat{p}_x . Hence we cannot simultaneously assign definite values to x and p_x , in agreement with the uncertainty principle. Since \hat{x} and \hat{H} do not commute, we cannot expect to assign definite values to the energy and x coordinate at the same time . A stationary state (which has a definite energy) shows a spread of possible values for x , the probabilities for observing various values of x being given by Born postulate.

For a state function ψ that is not an eigen function of \hat{A} , we get various possible outcomes when we measure A in identical system. We want some measure of the of the spread or dispersion in the set of observed values A_i . If $\langle A \rangle$ is the average of these values, then the deviation of each measurement from the average is $A_i - \langle A \rangle$. If we averaged all the deviations , we would get zero, since positive and negative deviations would cancel. Hence to make all